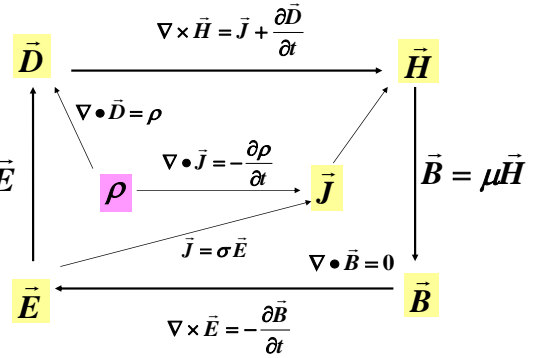


# Electromagnetic (EM) Waves

## TIME HARMONIC FORM OF MAXWELL'S EQUATIONS:

$$\begin{aligned}
 \nabla \cdot \vec{E} &= \rho_v / \epsilon \\
 \nabla \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} \\
 \nabla \cdot \vec{H} &= 0 \\
 \nabla \times \vec{H} &= \vec{J} + \frac{\partial \vec{D}}{\partial t}
 \end{aligned}
 \quad \left\{ \begin{array}{l} \vec{D} = \epsilon \vec{E} \\ \vec{B} = \mu \vec{H} \\ \vec{J} = \sigma \vec{E} \end{array} \right. \Rightarrow \left\{ \begin{array}{l} \nabla \cdot \vec{E} = \tilde{\rho}_v / \epsilon \\ \nabla \times \vec{E} = -j\omega \mu \vec{H} \\ \nabla \cdot \vec{H} = 0 \\ \nabla \times \vec{H} = j\omega(\epsilon - j\frac{\sigma}{\omega})\vec{E} = j\omega\epsilon_c \vec{E} \end{array} \right.$$

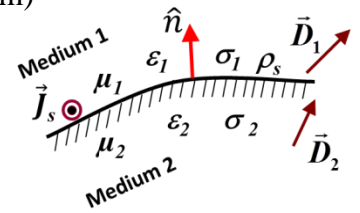


$$\vec{E}(x, y, z, t) = \Re\{ \tilde{E}(x, y, z) e^{j\omega t} \} \quad \epsilon_c = \epsilon - j\frac{\sigma}{\omega} = \epsilon' - j\epsilon''$$

CONSTANTS:  $\mu_0 \approx 4\pi \times 10^{-7} \text{ H/m}$     $\epsilon_0 \approx \frac{1}{36\pi} \times 10^{-9} \text{ F/m}$     $c = 3 \times 10^8 \text{ m/s}$

BOUNDARY CONDITIONS (Hold for both general time variation and phasor form)

$E_{1t} = E_{2t} \quad \text{or} \quad \hat{n} \times (\vec{E}_1 - \vec{E}_2) = 0$	$-H_{1t} + H_{2t} = J_s \quad \text{or} \quad \hat{n} \times (\vec{H}_1 - \vec{H}_2) = \vec{J}_s$
$D_{1n} - D_{2n} = \rho_s \quad \text{or} \quad \hat{n} \cdot (\vec{D}_1 - \vec{D}_2) = \rho_s$	$B_{1n} = B_{2n} \quad \text{or} \quad \hat{n} \cdot (\vec{B}_1 - \vec{B}_2) = 0$



PLANE WAVE IN SOURCE-FREE REGION (traveling in z-direction)

$$\vec{E}(z) = \hat{x}\tilde{E}_x(z) + \hat{y}\tilde{E}_y(z) \quad \tilde{E}_x(z) = \tilde{E}_x^+(z) + \tilde{E}_x^-(z) = E_{x0}^+ e^{-(\alpha+j\beta)z} + E_{x0}^- e^{(\alpha+j\beta)z}$$

$$\tilde{H}(z) = \hat{x}\tilde{H}_x(z) + \hat{y}\tilde{H}_y(z) \quad \gamma = \alpha + j\beta \quad \gamma^2 = -\omega^2 \mu \epsilon_c = -\omega^2 \mu (\epsilon' - j\epsilon'')$$

Speed:  $u_p = \frac{\omega}{\beta} = f\lambda = \frac{1}{\sqrt{\mu\epsilon}} = \frac{c}{\sqrt{\epsilon_r}}$ ;   Wavelength:  $\lambda = \frac{2\pi}{\beta} = \frac{2\pi}{k} = \frac{\lambda_0}{\sqrt{\epsilon_r}} = \frac{\lambda_0}{n}$ ;   Skin Depth:  $\delta_s = \frac{1}{\alpha}$

$$\vec{E} = -\eta_c \hat{z} \times \vec{H} \quad \vec{H} = \frac{1}{\eta_c} \hat{z} \times \vec{E} \quad \eta_c = \sqrt{\frac{\mu}{\epsilon_c}} = \sqrt{\frac{\mu}{\epsilon'} \left(1 - j\frac{\epsilon''}{\epsilon'}\right)^{-1/2}} \quad \eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} = 120\pi \Omega$$

## MATERIAL PROPERTIES:

	Any Medium	Lossless Medium ( $\sigma = 0$ )	Low-loss Medium ( $\epsilon''/\epsilon' \ll 1$ )	Good Conductor ( $\epsilon''/\epsilon' \gg 1$ )	Units
$\alpha =$	$\omega \left[ \frac{\mu\epsilon'}{2} \left[ \sqrt{1 + \left(\frac{\epsilon''}{\epsilon'}\right)^2} - 1 \right] \right]^{1/2}$	0	$\frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}}$	$\sqrt{\pi f \mu \sigma}$	(Np/m)
$\beta =$	$\omega \left[ \frac{\mu\epsilon'}{2} \left[ \sqrt{1 + \left(\frac{\epsilon''}{\epsilon'}\right)^2} + 1 \right] \right]^{1/2}$	$\omega \sqrt{\mu\epsilon}$	$\omega \sqrt{\mu\epsilon}$	$\sqrt{\pi f \mu \sigma}$	(rad/m)
$\eta_c =$	$\sqrt{\frac{\mu}{\epsilon'}} \left(1 - j\frac{\epsilon''}{\epsilon'}\right)^{-1/2}$	$\sqrt{\frac{\mu}{\epsilon}}$	$\sqrt{\frac{\mu}{\epsilon}}$	$(1 + j)\frac{\alpha}{\sigma}$	( $\Omega$ )
$u_p =$	$\omega/\beta$	$1/\sqrt{\mu\epsilon}$	$1/\sqrt{\mu\epsilon}$	$\sqrt{4\pi f/\mu\sigma}$	(m/s)
$\lambda =$	$2\pi/\beta = u_p/f$	$u_p/f$	$u_p/f$	$u_p/f$	(m)

Notes:  $\epsilon' = \epsilon$ ;  $\epsilon'' = \sigma/\omega$ ; in free space,  $\epsilon = \epsilon_0$ ,  $\mu = \mu_0$ ; in practice, a low-loss medium if  $\epsilon''/\epsilon' = \sigma/\omega\epsilon < 0.01$  and a good conducting medium if  $\epsilon''/\epsilon' > 100$ .

**Poynting vector:**  $\vec{S} = \vec{E} \times \vec{H}$  [W/m<sup>2</sup>];

Average power density:  $\vec{S}_{av} = \frac{1}{2} \text{Re}\{\vec{E} \times \vec{H}^*\} = \hat{z} \frac{|\vec{E}|^2}{2|\eta_c|} \cos \theta_{\eta_c}$  [W/m<sup>2</sup>]

Lossless media: Wavenumber  $k = \beta = \omega\sqrt{\mu\epsilon}$

**POLARIZATION:**  $\vec{E}(z) = (\hat{x}a_x + \hat{y}a_y e^{j\delta})e^{-jkz}$ ;  $\vec{E}(z, t) = \hat{x}a_x \cos(\omega t - \kappa z) + \hat{y}a_y \cos(\omega t - \kappa z + \delta)$

*Linear:*  $\delta=0$  or  $\pi$ ; *Circular:*  $a_x=a_y$  &  $\delta = \pm\pi/2$  (LHC/RHC); *Elliptical:* General

**NORMAL INCIDENCE:**

$$\begin{cases} \Gamma = \frac{E_o^r}{E_o^i} = \frac{\eta_{c2} - \eta_{c1}}{\eta_{c2} + \eta_{c1}} \\ \tau = \frac{E_o^t}{E_o^i} = \frac{2\eta_{c2}}{\eta_{c2} + \eta_{c1}} = 1 + \Gamma \end{cases} \quad \begin{cases} \vec{S}_{av1}(z) = \hat{z} \frac{|E_o^i|^2}{2|\eta_c|} [(e^{-2\alpha z} - |\Gamma|^2 e^{2\alpha z}) \cos \theta_{\eta_c} - 2|\Gamma| \sin(2\beta z + \theta_\Gamma) \sin \theta_{\eta_c}] \\ \vec{S}_{av2}(z) = \hat{z} \frac{|E_o^i|^2}{2|\eta_{c2}|} |\tau|^2 e^{-2\alpha_2 z} \cos \theta_{\eta_{c2}} \end{cases}$$

$\vec{E}_i$ & $\vec{E}_r$	$\vec{E}_t$
$\vec{H}_i$ & $\vec{H}_r$	$\vec{H}_t$
$(\mu_1, \epsilon_1, \sigma_1)$	$(\mu_2, \epsilon_2, \sigma_2)$
$(\gamma_1, \eta_{c1})$	$(\gamma_2, \eta_{c2})$

For x-polarized EM wave traveling in z-direction:

$$\begin{cases} \vec{E}_1(z) = \vec{E}^i + \vec{E}^r = \hat{x}E_o^i(e^{-\gamma_1 z} + \Gamma e^{\gamma_1 z}) \\ \vec{H}_1(z) = \vec{H}^i + \vec{H}^r = \hat{y} \frac{E_o^i}{\eta_{c1}}(e^{-\gamma_1 z} - \Gamma e^{\gamma_1 z}) \end{cases} \quad \begin{cases} \vec{E}_t(z) = \hat{x}\tau E_o^i e^{-\gamma_2 z} \\ \vec{H}_t(z) = \hat{y}\tau \frac{E_o^i}{\eta_{c2}} e^{-\gamma_2 z} \end{cases}$$

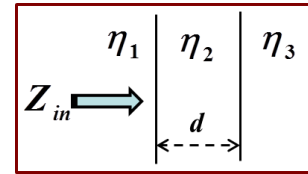
Multiple Boundaries:

Wave impedance

$$Z(z) = \frac{\vec{E}(z)}{\vec{H}(z)} = \eta \frac{1 + \Gamma(z)}{1 - \Gamma(z)}$$

Input impedance

$$Z_{in} = \eta_{in} = \eta_2 \frac{\eta_3 + j\eta_2 \tan \beta_2 d}{\eta_2 + j\eta_3 \tan \beta_2 d}$$



**OBLIQUE INCIDENCE:**

Arbitrary ( $\vec{k}$ ) Direction:  $\vec{E} = \vec{E}_m e^{-j\vec{k}\cdot\vec{r}} = \hat{a}_E E_m e^{-jkx_i}$      $\vec{H} = \frac{1}{\eta} \hat{k} \times \vec{E} = \hat{a}_H \frac{E_m}{\eta} e^{-jkx_i}$

e.g., for the figure below,  $\vec{E}_\perp^i = \hat{y}E_{\perp o}^i e^{-jk_i x_i}$ ,  $\vec{H}_\perp^i = \hat{y}_i \frac{E_{\perp o}^i}{\eta_i} e^{-jk_i x_i}$ ,

where  $x_i = \hat{k}_i \cdot \vec{r} = x \sin \theta_i + z \cos \theta_i$ ,  $\hat{y}_i = -\hat{x} \cos \theta_i + \hat{z} \sin \theta_i$ .

**Snell's Law:**  $\theta_i = \theta_r \equiv \theta_t$

Index of Refraction:  $n = \sqrt{\epsilon_r}$  (lossless non-magnetic)

$$k_1 \sin \theta_i = k_2 \sin \theta_t \quad \text{or} \quad n_1 \sin \theta_i = n_2 \sin \theta_t$$

**Perpendicular Polarization:**

$$\Gamma_\perp = \frac{E_{\perp o}^r}{E_{\perp o}^i} = \frac{\eta_2 \cos \theta_i - \eta_1 \cos \theta_t}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t}$$

$$\tau_\perp = \frac{E_{\perp o}^t}{E_{\perp o}^i} = \frac{2\eta_2 \cos \theta_i}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t} = 1 + \Gamma_\perp$$

**Parallel Polarization:**

$$\Gamma_\parallel = \frac{E_{\parallel o}^r}{E_{\parallel o}^i} = \frac{\eta_2 \cos \theta_t - \eta_1 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i}$$

$$\tau_\parallel = \frac{E_{\parallel o}^t}{E_{\parallel o}^i} = \frac{2\eta_2 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i} = (1 + \Gamma_\parallel) \frac{\cos \theta_i}{\cos \theta_t}$$

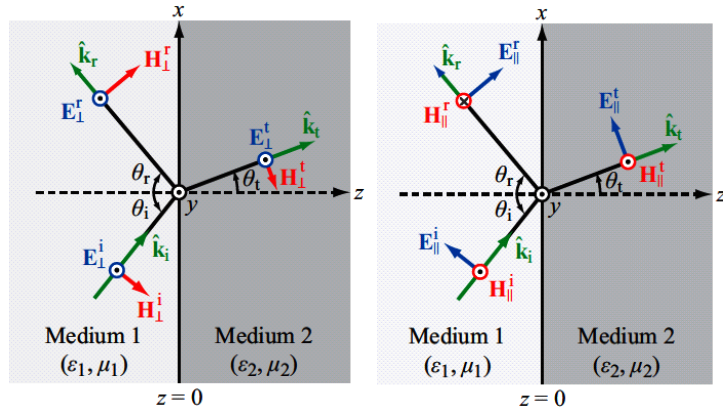
**Power Reflectivity and Transmissivity:**

$$R = |\Gamma|^2 \quad T = 1 - R$$

**Critical Angle:**  $\sin \theta_c = \frac{n_2}{n_1} = \sqrt{\frac{\epsilon_2}{\epsilon_1}}$

**Brewster Angle:**  $\theta_{B\parallel} = \sin^{-1} \frac{1}{\sqrt{1 + \epsilon_1 / \epsilon_2}}$   
(for  $\mu_1 = \mu_2$ )

$$= \tan^{-1} \sqrt{\frac{\epsilon_2}{\epsilon_1}}$$



(a) Perpendicular polarization

(b) Parallel polarization